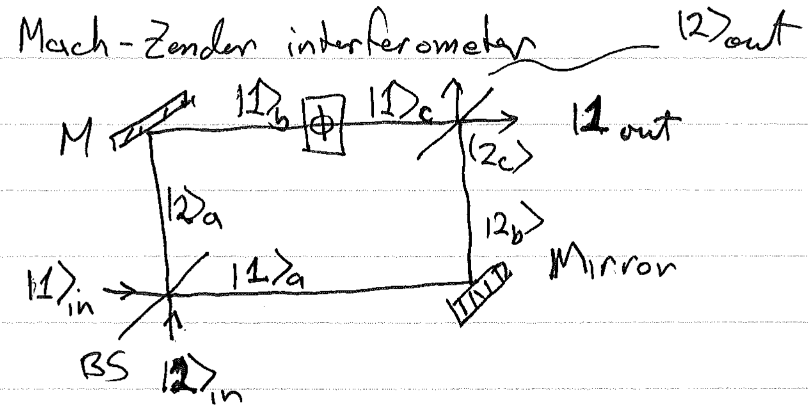
Physics 566 - Quantum Optics

Problem Set #1: Solutions

Problem 1: Equivalence of a Mach-Zender and SU(2) interferometer

A Mach-Zender interferometer



(a) We want to show that the MZ transformation on the two modes $|1\rangle$ and $|2\rangle$ is equivalent to the SU(2) transformation on a 2D Hilbert space, as given in the assignment

Recall:
$$\begin{cases} \frac{\pi}{2}\text{ x-rotation: } & e^{-i\frac{\pi}{2}\frac{\sigma_x}{2}} \\ \pi \text{ x-rotation: } & e^{-i\pi\sigma_x/2} \\ \phi \text{ z-rotation } & e^{-i\phi\,\sigma_z/2} \end{cases}$$

Sequence: Remember, first operation on the right and then sequentially multiply on left

$$U_{total} = e^{-i\frac{\pi}{4}\hat{\sigma}_x}\, e^{-i\phi\frac{\hat{\sigma}_z}{2}}\, e^{-i\frac{\pi}{2}\hat{\sigma}_x}\, e^{-i\frac{\pi}{4}\hat{\sigma}_x}$$

$$= e^{-i\frac{\pi}{4}\hat{\sigma}_x}\, e^{-i\phi\frac{\hat{\sigma}_z}{2}}\, e^{-i\frac{3\pi}{4}\hat{\sigma}_x}$$

$$\Rightarrow U_{total} = \underbrace{e^{-i\frac{\pi}{4}\hat{\sigma}_x}\, e^{-i\phi\frac{\hat{\sigma}_z}{2}} e^{+i\frac{\pi}{4}\hat{\sigma}_x}}_{= \; e^{+i\phi\frac{\hat{\sigma}_y}{2}}} \quad \underbrace{e^{i\pi\sigma_x}}_{= \; -\hat{\mathbb{1}}}$$

negligible overall phase

This follows since $D_x(\frac{\pi}{2}) = e^{-i\frac{\pi}{4}\hat{\sigma}_x}$

rotates 90° around x-axis and

$$D_x^{\dagger}(\tfrac{\pi}{2})\,\hat{\sigma}\, D_x(\tfrac{\pi}{2}) = \hat{\sigma}_y$$

$$\therefore U_{total} \equiv e^{i\phi\frac{\hat{\sigma}_y}{2}} = \cos\frac{\phi}{2}\,\mathbb{1} + i\sin\frac{\phi}{2}\,\hat{\sigma}_y$$

$$= \cos\frac{\phi}{2}\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} + i\sin\frac{\phi}{2}\begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$$

$$\boxed{U_{total} = \begin{bmatrix} \cos\frac{\phi}{2} & \sin\frac{\phi}{2} \\ -\sin\frac{\phi}{2} & \cos\frac{\phi}{2} \end{bmatrix}}$$

Does this achieve the same transformation as the MZ interferometer?

50-50 Beam splitter $\Leftrightarrow$ $\frac{\pi}{2}$-X rotation?

$$e^{-i\frac{\pi}{4}\sigma_x} = \cos\frac{\pi}{4}\,\hat{\mathbb{1}} - i\sin\frac{\pi}{4}\,\sigma_x$$

$$= \frac{1}{\sqrt{2}}\begin{bmatrix} 1 & -i \\ -i & 1 \end{bmatrix}$$

$$e^{-i\frac{\pi}{4}\sigma_x}|1_{in}\rangle = \frac{1}{\sqrt{2}}|1\rangle_a - \frac{i}{\sqrt{2}}|2\rangle_a$$

$$P_{1a} = \left|\frac{1}{\sqrt{2}}\right|^2 = P_{2a} = \left|-\frac{i}{\sqrt{2}}\right|^2 = \frac{1}{2}$$ ✓

- 50-50 split, but with a particular phase ✓

<u>Mirrors $\Leftrightarrow$ π-x-rotation?</u>

$$e^{-i\frac{\pi}{2}\sigma_x} = \cos\frac{\pi}{2}\,\mathbb{1} - i\sin\frac{\pi}{2}\,\sigma_x = -i\,\hat{\sigma}_x$$

$$= -i\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

$$\left.\begin{aligned} e^{-i\frac{\pi}{2}\sigma_x}|1_a\rangle &= -i\,|2_b\rangle \\ e^{-i\frac{\pi}{2}\sigma_x}|2_a\rangle &= -i\,|1_a\rangle \end{aligned}\right\}$$ Flips modes (up to overall negligible phase) ✓

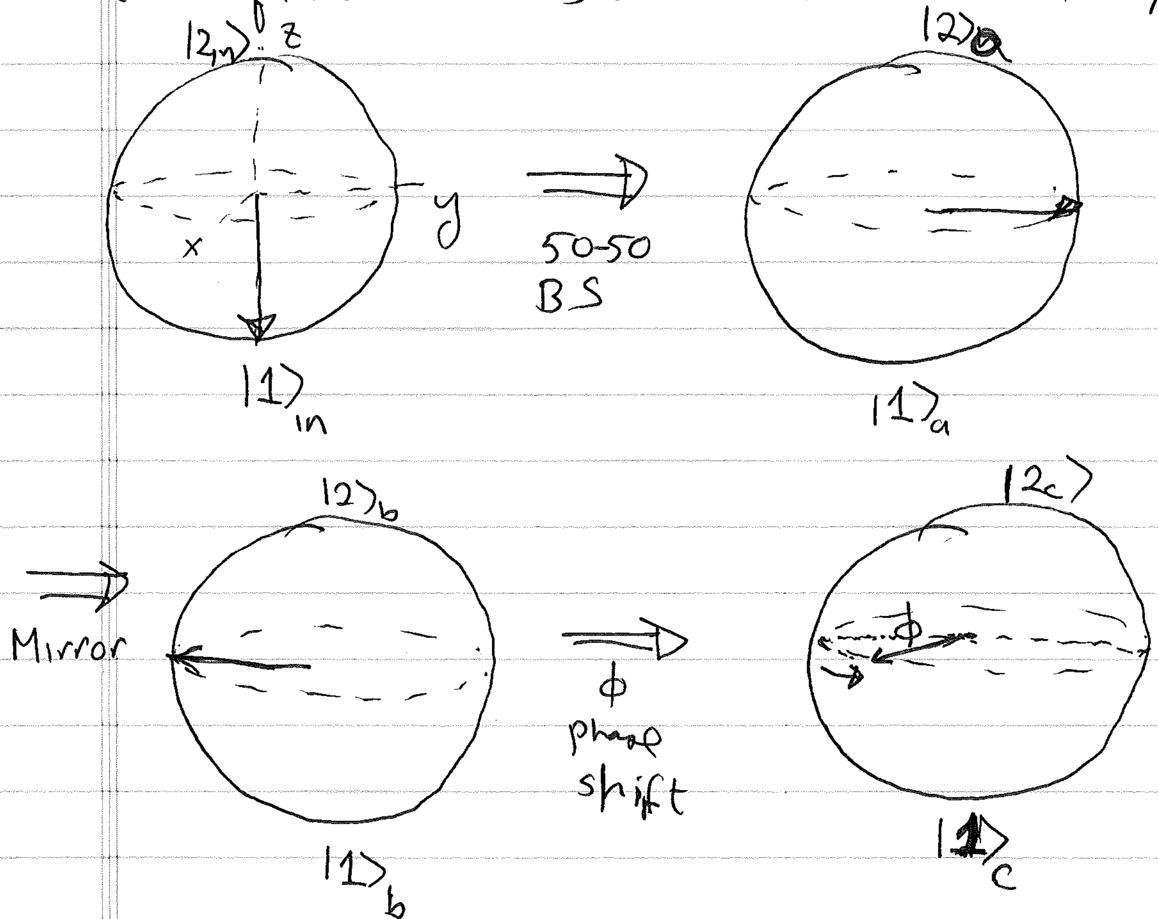
<u>Phase shift $\Leftrightarrow$ ϕ-z-rotation?</u>

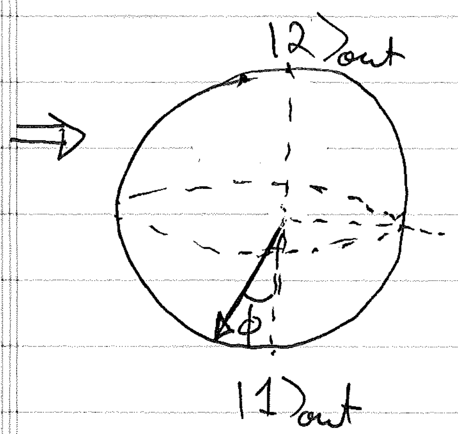
$$e^{-i\phi\frac{\hat{\sigma}_z}{2}} = \cos\frac{\phi}{2}\,\mathbb{1} - i\sin\frac{\phi}{2}\,\hat{\sigma}_z$$

$$= \begin{bmatrix} \cos\frac{\phi}{2} - i\sin\frac{\phi}{2} & 0 \\ 0 & \cos\frac{\phi}{2} + i\sin\frac{\phi}{2} \end{bmatrix} = \begin{bmatrix} e^{-i\phi/2} & 0 \\ 0 & e^{i\phi/2} \end{bmatrix}$$

$$\Rightarrow \left.\begin{aligned} e^{-i\phi\frac{\hat{\sigma}_z}{2}}|1_b\rangle &= e^{-i\phi/2}|1_c\rangle \\ e^{-i\phi\frac{\hat{\sigma}_z}{2}}|2_b\rangle &= e^{+i\phi/2}|2_c\rangle \end{aligned}\right\}$$ relative phase of ϕ ✓

(b) Sequence of transformations on Bloch Sphere





<u>Note</u>: When $\phi = 0$
(a balanced interferometer)
$|1\rangle_{in} \Rightarrow |1\rangle_{out}$
and zero probability in $|2\rangle_{out}$

(c) Given $U_{total} = \begin{bmatrix} \cos\frac{\phi}{2} & \sin\frac{\phi}{2} \\ -\sin\frac{\phi}{2} & \cos\frac{\phi}{2} \end{bmatrix}$

$$P_{1out} = |\langle 1_{out} | U_{total} | 1_{in} \rangle|^2 = \cos^2\frac{\phi}{2}$$ ✓

(d) Now consider mixed states, described in general by the density operator, $\hat{\rho}_{in}$.

Then, $\hat{\rho}_{out} = \hat{U}_{total} \hat{\rho}_{in} \hat{U}^{\dagger}_{total}$

Given $\hat{\rho}_{in} = P_{1in} |1_{in}\rangle\langle 1_{in}| + P_{2in} |2_{in}\rangle\langle 2_{in}|$

Then $P_{1out} = P_{1in} |\langle 1_{out} | U_{total} | 1_{in}\rangle|^2$

$$+ P_{2in} |\langle 2_{out} | U_{total} | 1_{in}\rangle|^2$$

Aside: $|\langle 1_{out} | U_{total} | 1_{in}\rangle|^2 = \cos^2\frac{\phi}{2}$

$$|\langle 2_{out} | U_{total} | 1_{in}\rangle|^2 = \sin^2\frac{\phi}{2}$$

$\Rightarrow$ (i) With $\hat{\rho}_{in} = \frac{1}{2}|1_{in}\rangle\langle 1_{in}| + \frac{1}{2}|2_{in}\rangle\langle 2_{in}|$

$$\boxed{P_{1out} = \frac{1}{2}\cos^2\frac{\phi}{2} + \frac{1}{2}\sin^2\frac{\phi}{2} = \frac{1}{2}}$$

(ii) With $\hat{\rho}_{in} = \frac{1}{3}|1_{in}\rangle\langle 1_{in}| + \frac{2}{3}|2_{in}\rangle\langle 2_{in}|$

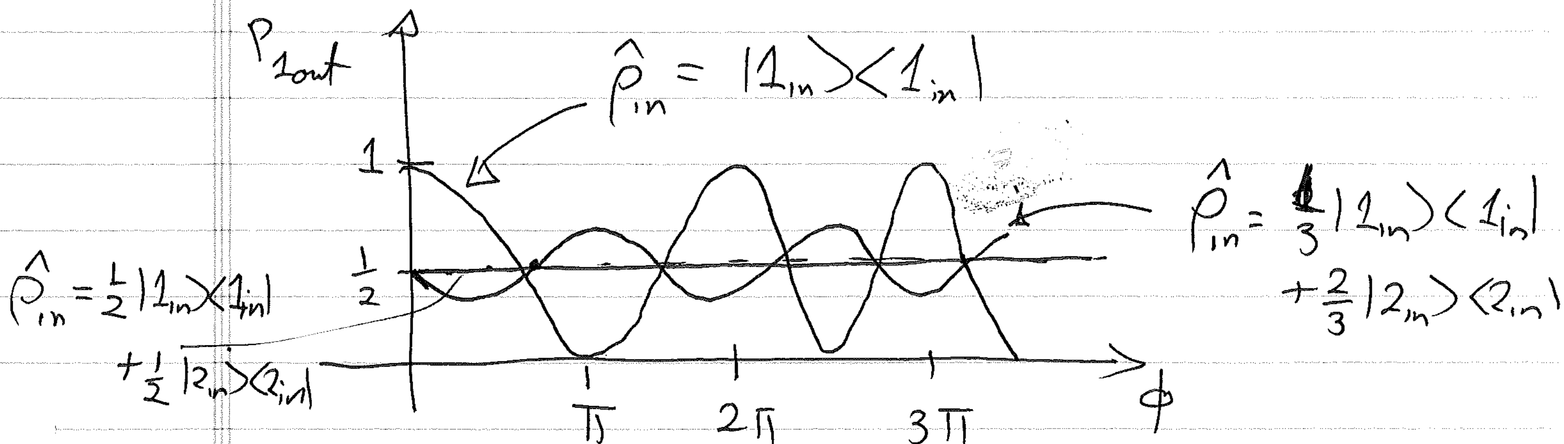
$$P_{1out} = \frac{1}{3}\cos^2\frac{\phi}{2} + \frac{2}{3}\sin^2\frac{\phi}{2} =$$

$\Rightarrow$ (d ii)

$$P_{1out} = \frac{1}{3}\left(\frac{1+\cos\phi}{2}\right) + \frac{2}{3}\left(\frac{1-\cos\phi}{2}\right)$$

$$\boxed{P_{1out} = \frac{1}{2} - \frac{1}{6}\cos\phi}$$

Sketch



- For the pure state we have maximum visibility
- For the completely mixed state, there is no coherence, and no interference
- For the partially mixed state, there is a non-zero, but not perfect visibility.

Problem 2

(a) $\mathrm{Tr}(|\phi\rangle\langle\xi|) = \sum_n \langle n|\phi\rangle\langle\xi|n\rangle$ for basis $\{n\}$

$$= \sum_n \langle\xi|n\rangle\langle n|\phi\rangle = \langle\xi|\underbrace{\left(\sum_n |n\rangle\langle n|\right)}_{=\mathbb{1}}|\phi\rangle$$

$$= \langle\xi|\phi\rangle$$

Trace turns outer product to inner product

(a) Consider a statistical mixture

$$\hat{\rho} = P_+ |+_z\rangle\langle +_z| + P_- |-_z\rangle\langle -_z|$$

where $P_\pm = \frac{1}{2}\left(1 \pm \frac{1}{\sqrt{2}}\right)$

Density matrix $\rho_{ij} = \langle i|\hat{\rho}|j\rangle$

In basis $|\pm_z\rangle$ $\hat{\rho} \underset{z}{\doteq} \frac{1}{2}\begin{pmatrix} 1+\frac{1}{\sqrt{2}} & 0 \\ 0 & 1-\frac{1}{\sqrt{2}} \end{pmatrix}$

In basis $\left(|\pm_x\rangle = \frac{1}{\sqrt{2}}(|+_z\rangle \pm |-_z\rangle)\right)$

$$\hat{\rho} \underset{x}{\doteq} \begin{pmatrix} \langle +_x|\hat{\rho}|+_x\rangle & \langle +_x|\hat{\rho}|-_x\rangle \\ \langle -_x|\hat{\rho}|+_x\rangle & \langle -_x|\hat{\rho}|-_x\rangle \end{pmatrix} = \frac{1}{2}\begin{pmatrix} 1 & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & 1 \end{pmatrix}$$

Note: In this basis the density operator has off-diagonal elements. Nonetheless, it is a mixed state:

$$\mathrm{Tr}(\hat{\rho}^2) = \frac{3}{4}$$

The Bloch vector can be seen immediately from the form in the z-basis.

$$\mathrm{Tr}(\hat{\rho}\,\hat{\sigma}_x) = \mathrm{Tr}(\hat{\rho}\,\hat{\sigma}_y) = 0$$

$$\mathrm{Tr}(\hat{\rho}\,\hat{\sigma}_z) = P_+ - P_- = \frac{1}{\sqrt{2}}$$

$$\Rightarrow \boxed{\vec{Q} = \frac{1}{\sqrt{2}}\vec{e}_z} \qquad \text{mixed state } |\vec{Q}| < 1$$

(c) Now we have a state

$$\hat{\rho} = \frac{1}{2}|+_{n_1}\rangle\langle +_{n_1}| + \frac{1}{2}|+_{n_2}\rangle\langle +_{n_2}|$$

where $|+_n\rangle\langle +_n| = \frac{1}{2}(\hat{\mathbb{1}} + \vec{e}_n \cdot \hat{\vec{\sigma}})$ from Prob 1

$$\vec{e}_{n_{1/2}} = \frac{1}{\sqrt{2}}(\vec{e}_z \pm \vec{e}_x)$$

$$\Rightarrow \hat{\rho} = \frac{1}{2}\hat{\mathbb{1}} + \frac{1}{4}(\vec{e}_{n_1} + \vec{e}_{n_2}) \cdot \hat{\vec{\sigma}}$$

$$= \frac{1}{2}\hat{\mathbb{1}} + \frac{1}{4}\left(\frac{2}{\sqrt{2}}\vec{e}_z\right) \cdot \hat{\vec{\sigma}}$$

$$= \frac{1}{2}\left(\hat{\mathbb{1}} + \frac{1}{\sqrt{2}}\vec{e}_z\right) \cdot \vec{\sigma} \underset{z}{=} \begin{bmatrix} \frac{1}{2}(1+\frac{1}{\sqrt{2}}) & 0 \\ 0 & \frac{1}{2}(1-\frac{1}{\sqrt{2}}) \end{bmatrix}$$

Same as $\hat{\rho}$ in part (b)!

Moral of the story: The ensemble decomposition is not unique. In fact, we can take any density matrix for a two-level system, described uniquely in terms of its Bloch vector $\vec{Q}$ and decompose it in terms of an ensemble of any two pure states described by unit vector $\vec{e}_n$ with probability p_n if $\vec{Q} = p_{n_1}\vec{e}_{n_1} + p_{n_2}\vec{e}_{n_2}$.

(c) Two statistical mixtures

$$\hat{\rho}_1 = \sum_n p_n |+_n\rangle\langle +_n|$$

$$\hat{\rho}_2 = \sum_m q_m |+_m\rangle\langle +_m|$$

Aside: $\underbrace{|+_n\rangle\langle +_n|}_{\text{Projector}} = \frac{1}{2}(\hat{\mathbb{1}} + \hat{\sigma}_n)$ $\left(\begin{array}{l}\text{where} \\ \hat{\sigma}_n = \vec{e}_n \cdot \hat{\vec{\sigma}}\end{array}\right)$

$$\Rightarrow \hat{\rho}_1 = \underbrace{\left(\sum_n p_n\right)}_{=1} \frac{1}{2}\hat{\mathbb{1}} + \frac{1}{2}\underbrace{\left(\sum_n p_n \vec{e}_n\right)}_{\vec{Q}_1} \cdot \hat{\vec{\sigma}}$$

$$\Rightarrow \hat{\rho}_1 = \frac{1}{2}\hat{\mathbb{1}} + \frac{1}{2}\vec{Q}_1 \cdot \hat{\vec{\sigma}}$$

Similarly $\hat{\rho}_2 = \frac{1}{2}\hat{\mathbb{1}} + \frac{1}{2}\vec{Q}_2 \cdot \hat{\vec{\sigma}}$

where $\vec{Q}_2 = \sum_m q_m \vec{e}_m$

Thus $\hat{\rho}_1 = \hat{\rho}_2 \iff \vec{Q}_1 = \vec{Q}_2$

Problem 3: Ambiguity of ensemble decomposition

Let $\hat{\rho}_1 = \sum_i p_i |\psi_i\rangle\langle\psi_i|$, $\hat{\rho}_2 = \sum_j q_j |\phi_j\rangle\langle\phi_j|$

Prove $\hat{\rho}_1 = \hat{\rho}_2$ iff $\sqrt{q_j}\,|\phi_j\rangle = \sum_i U_{ji} \sqrt{p_i}\,|\psi_i\rangle$

where U_{ji} are elements of unitary matrix.

Proof:

For convenience, define $|\bar{\phi}_j\rangle \equiv \sqrt{q_j}\,|\phi_j\rangle$

$|\bar{\psi}_i\rangle \equiv \sqrt{p_i}\,|\psi_i\rangle$

so $\langle\bar{\phi}_j|\bar{\phi}_j\rangle = q_j$ $\langle\bar{\psi}_i|\bar{\psi}_i\rangle = p_i$

(i) Assume $|\bar{\phi}_j\rangle = \sum_i U_{ji} |\bar{\psi}_i\rangle$ U_{ji} elements of unitary matrix

Consider $\hat{\rho}_2 = \sum_j |\bar{\phi}_j\rangle\langle\bar{\phi}_j| = \sum_{jk} U^*_{jk} U_{ji} |\bar{\psi}_i\rangle\langle\bar{\psi}_k|$

Aside: $(U_{jk})^* = U^\dagger_{kj}$

$$\Rightarrow \hat{\rho}_2 = \sum_{ik} \underbrace{\left(\sum_j U^\dagger_{kj} U_{ji}\right)}_{\delta_{ik}} |\bar{\psi}_i\rangle\langle\bar{\psi}_k|$$

$$\Rightarrow \hat{\rho}_2 = \sum_i |\bar{\psi}_i\rangle\langle\bar{\psi}_i| = \hat{\rho}_1 \checkmark$$

(ii) Now assume $\hat{\rho}_1 = \hat{\rho}_2 \equiv \hat{\rho}$

$\hat{\rho}$ being a Hermitian operator can be diagonalized

$$\Rightarrow \quad \hat{\rho} = \sum_\alpha \lambda_\alpha |e_\alpha\rangle\langle e_\alpha|$$

$$\text{where } \begin{cases} \sum_\alpha \lambda_\alpha = 1 & \text{with } \lambda_\alpha \text{ real}, \quad 0 \le \lambda_\alpha \le 1 \\ \langle e_\alpha | e_\beta \rangle = \delta_{\alpha\beta} \end{cases}$$

Let $|\bar{e}_\alpha\rangle = \sqrt{\lambda_\alpha}\, |e_\alpha\rangle \quad \Rightarrow \quad \hat{\rho} = \sum_\alpha |\bar{e}_\alpha\rangle\langle \bar{e}_\alpha|$

$$\Rightarrow \quad \sum_i |\Psi_i\rangle\langle\Psi_i| = \sum_\alpha |\bar{e}_\alpha\rangle\langle\bar{e}_\alpha| = \sum_j |\Phi_j\rangle\langle\Phi_j|$$

we seek the relationship between $\{|\Psi_i\rangle\}$ and $\{|\Phi_j\rangle\}$

First note $\{|e_\alpha\rangle\}$ form a basis for the Hilbert space (with $\lambda_\alpha = 0$ for those vectors not in $\hat{\rho}$)

$$\Rightarrow \quad |\Psi_i\rangle = \sum_\alpha |e_\alpha\rangle\langle e_\alpha|\Psi_i\rangle = \sum_\alpha |\bar{e}_\alpha\rangle \frac{\langle e_\alpha|\Psi_i\rangle}{\sqrt{\lambda_\alpha}}$$

$$= \sum_\alpha M_{i\alpha} |\bar{e}_\alpha\rangle$$

$$\text{where} \quad M_{i\alpha} = \frac{\langle e_\alpha|\Psi_i\rangle}{\sqrt{\lambda_\alpha}}$$

Now: $\sum_i M_{i\alpha} M^*_{i\beta} = \sum_i \frac{\langle e_\alpha | \Psi_i \rangle \langle \Psi_i | e_\beta \rangle}{\sqrt{\lambda_\alpha \lambda_\beta}}$

$$= \frac{1}{\sqrt{\lambda_\alpha \lambda_\beta}} \langle e_\alpha | \left(\sum_i |\Psi_i\rangle\langle\Psi_i| \right) | e_\beta \rangle$$

$$= \frac{1}{\sqrt{\lambda_\alpha \lambda_\beta}} \langle e_\alpha | \hat{\rho} | e_\beta \rangle = \frac{\lambda_\alpha \delta_{\alpha\beta}}{\sqrt{\lambda_\alpha \lambda_\beta}} = \delta_{\alpha\beta}$$

$\Rightarrow$ When arranged in a matrix, the columns of $M_{i\alpha}$ are orthonormal

(Subtle point: $M_{i\alpha}$ need not be square here, since # of pure states in the $\{|\Psi_i\rangle\}$ need not be the dimension of Hilbert space. However, we can always append extra columns in the orthogonal space to make $M_{i\alpha}$ unitary.)

Thus since $|\Psi_i\rangle = \sum_\alpha M_{i\alpha} |\bar{e}_\alpha\rangle$

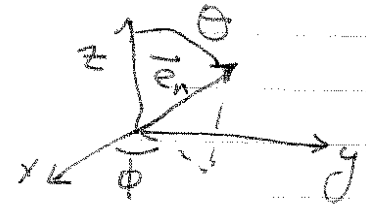
$$|\Phi_j\rangle = \sum_\beta N_{j\beta} |\bar{e}_\beta\rangle$$

$$|\Phi_j\rangle = \sum_i U_{ji} |\Psi_i\rangle$$

where $U = NM^\dagger$ q.e.d.

Problem 3

Given unit direction in space $\mathbb{R}^3$



$$\vec{e}_n = \cos\theta\,\vec{e}_z + \sin\theta(\cos\phi\,\vec{e}_x + \sin\phi\,\vec{e}_y)$$

(a)

- We showed in class, spin up along $\vec{e}_n$ is

$$|+_n\rangle = \cos\tfrac{\theta}{2}|+_z\rangle + e^{i\phi}\sin\tfrac{\theta}{2}|-_z\rangle$$

- An arbitrary pure state for two levels

$$|\psi\rangle = \alpha|+_z\rangle + \beta|-_z\rangle \qquad \underline{|\alpha|^2+|\beta|^2=1}$$

$$= |\alpha|\,|+_z\rangle + e^{i\delta_{\alpha\beta}}\,|\beta|\,|-_z\rangle$$

where $\delta_{\alpha\beta} = \mathrm{Arg}(\beta) - \mathrm{Arg}(\alpha)$

Thus $|\psi\rangle$ is of the form $|+_n\rangle$ with

$$\theta = 2\cos^{-1}(|\alpha|) \qquad \phi = \delta_{\alpha\beta}$$

q.e.d.

Note: This fact is <u>uniquely</u> <u>true</u> for spin 1/2

For $J > 1/2$, not all pure states can be mapped to directions in space.

(b) Consider a projector onto a pure state

$$\hat{\mathbb{P}}_n \equiv |+_n\rangle\langle +_n|$$

Every operator $\hat{A} = \frac{1}{2}\left(\mathrm{Tr}(\hat{A})\hat{\mathbb{1}} + \vec{A}\cdot\hat{\vec{\sigma}}\right)$

where $\vec{A} = \mathrm{Tr}(\hat{A}\,\hat{\vec{\sigma}})$

$$\mathrm{Tr}(\hat{\mathbb{P}}_n) = \langle +_n|+_n\rangle = 1$$

$\mathrm{Tr}(\hat{\mathbb{P}}_n\,\hat{\vec{\sigma}}) = \langle +_n|\hat{\vec{\sigma}}|+_n\rangle = \vec{Q}$ Bloch vector

But for a pure state $\vec{Q} = \vec{e}_n$

$\Rightarrow |+_n\rangle\langle +_n| = \frac{1}{2}\left(\hat{\mathbb{1}} + \vec{e}_n\cdot\hat{\vec{\sigma}}\right)$ q.e.d.

Explicitly $Q_x = \langle +_n|\hat{\sigma}_x|+_n\rangle = \langle +_n|+_z\rangle\langle -_z|+_n\rangle + c.c.$

$$= \left(\cos\tfrac{\theta}{2}\right)\left(e^{i\phi}\sin\tfrac{\theta}{2}\right) + c.c.$$

$$= \cos\phi\left(2\sin\tfrac{\theta}{2}\cos\tfrac{\theta}{2}\right) = \cos\phi\sin\theta$$

$$Q_y = \langle +_n|\hat{\sigma}_y|+_n\rangle = \frac{\langle +_n|+_z\rangle\langle -_z|+_n\rangle - c.c.}{i}$$

$$= \sin\phi\left(2\sin\tfrac{\theta}{2}\cos\tfrac{\theta}{2}\right) = \sin\phi\sin\theta$$

$$Q_z = \langle +_n|\hat{\sigma}_z|+_n\rangle = |\langle +_n|+_z\rangle|^2 - |\langle +_n|-_z\rangle|^2$$

$$= \cos^2\tfrac{\theta}{2} - \sin^2\tfrac{\theta}{2} = \cos\theta$$

$$\Rightarrow \vec{Q} = \sin\theta(\cos\phi\,\vec{e}_x + \sin\phi\,\vec{e}_y) + \cos\theta\,\vec{e}_z$$

$$= \vec{e}_n$$

(c) Consider

$$|\langle +_n | +_{n'} \rangle| = \sqrt{|\langle +_n | +_{n'} \rangle|^2}$$

Now $|\langle +_n | +_{n'} \rangle|^2 = \mathrm{Tr}\left(|+_n\rangle\langle +_n| \; |+_{n'}\rangle\langle +_{n'}| \right)$

$$= \mathrm{Tr}\left[\tfrac{1}{2}(\hat{\mathbb{1}} + \hat{\sigma}_n) \; \tfrac{1}{2}(\mathbb{1} + \hat{\sigma}_{n'}) \right]$$

$$= \frac{1}{4}\underbrace{\mathrm{Tr}(\hat{\mathbb{1}})}_{=2} + \frac{1}{2}\underbrace{\mathrm{Tr}(\hat{\sigma}_n + \hat{\sigma}_{n'})}_{\to 0} + \frac{1}{4}\mathrm{Tr}(\hat{\sigma}_n \hat{\sigma}_{n'})$$

<u>Aside:</u> $\mathrm{Tr}(\hat{\sigma}_n \hat{\sigma}_{n'}) = \mathrm{Tr}(\vec{e}_n \cdot \hat{\vec{\sigma}} \; \vec{e}_{n'} \cdot \vec{\sigma})$

$$= 2\, \vec{e}_n \cdot \vec{e}_{n'}$$

$$\Rightarrow |\langle +_n | +_{n'} \rangle| = \sqrt{\tfrac{1}{2}(1 + \vec{e}_n \cdot \vec{e}_{n'})} = \sqrt{\frac{1 + \cos\Theta}{2}}$$

where $\Theta = \cos^{-1}(\vec{e}_n \cdot \vec{e}_{n'})$

$$\Rightarrow \boxed{|\langle +_n | +_{n'} \rangle| = \left|\cos\left(\frac{\Theta}{2}\right)\right|}$$

Anti-podal states on Bloch sphere are orthogonal $\Theta = \pi$

(d) Poincaré sphere

if $|+_z\rangle \Rightarrow$ right-hand circular $= (\vec{e}_x + i\vec{e}_y)/\sqrt{2}$

$|-_z\rangle \Rightarrow$ left-hand circular $= (\vec{e}_x - i\vec{e}_y)/\sqrt{2}$

$$|\pm_x\rangle = \frac{|+_z\rangle \pm |-_z\rangle}{\sqrt{2}} = \begin{cases} \vec{e}_x \\ \vec{e}_y \end{cases}$$ linear — Horizontal / vertical

$$|\pm_y\rangle = \frac{|+_z\rangle \pm i|-_z\rangle}{\sqrt{2}} = \begin{cases} (\vec{e}_x + \vec{e}_y)/\sqrt{2} \\ (\vec{e}_x - \vec{e}_y)/\sqrt{2} \end{cases}$$ linear — 45° / -45°